

Formulazioni deboli

Integrale di Galerkin:

$$\int_R (Au_n - f) \xi_i dR = 0 \quad i = 1, 2, \dots, n$$

La 1^a identità di Green abbassa l'ordine delle derivate su u_n e l'alza su ξ_i . Si tratta dell'uso, anche ripetuto, dell'integrazioni per parti. Esempio con A operatore differenziale del 2° ordine:

$$\int_R (Au_n) \xi_i dR = \int_R u_n A^* \xi_i dR + \int_{\partial R} [G(u_n) F(\xi_i) - F(u_n) G(\xi_i)] d\Gamma$$

A^* : operatore aggiunto di A ($A^* = A$ se A è simmetrico)

F, G : operatori differenziali che saturano naturalmente della integrazione per parti

Esempi di formulazioni deboli: $A = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$

$$\int_R (Au) v dR = \int_R \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) v dR = - \int_R \left(\frac{\partial u}{\partial x} \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \frac{\partial v}{\partial y} \right) dR$$

$$+ \int_{\partial R} \frac{\partial u}{\partial n} v d\Gamma = \int_R \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) u dR + \int_{\partial R} \frac{\partial u}{\partial n} v d\Gamma - \int_{\partial R} u \frac{\partial v}{\partial n} d\Gamma$$

(2 applicazioni del lemma di Green)

In questo caso $A = A^*$ (perché Laplace è simmetrico)

Inoltre $F = 1$ e $G = G^* = \partial/\partial n$

Ad ogni integrazione per parti corrisponde una formulazione debole.

Se le condizioni al contorno sono trattate correttamente, quando $n \rightarrow \infty$, la soluzione debole $u_n \rightarrow u$, soluzione della PDE

Importanza ed utilità delle formulazioni deboli:

Esempio con l'eq. bi-armonica:

$$\Delta^4 u = f \quad \text{cioè}$$
$$\Delta^2(\Delta^2 u) - f = \frac{\partial^4 u}{\partial x^4} + 2 \frac{\partial^4 u}{\partial x^2 \partial y^2} + \frac{\partial^4 u}{\partial y^4} - f = 0$$

Integrale i-esimo di Galerkin:

$$\int_R (\Delta^4 u_n - f) \zeta_i \, dR = 0$$

1^a applicazione della formulazione debole:

$$\int_R \Delta^2(\Delta^2 u_n) \zeta_i \, dR = - \int_R \left[\frac{\partial}{\partial x} (\Delta^2 u_n) \frac{\partial \zeta_i}{\partial x} + \frac{\partial}{\partial y} (\Delta^2 u_n) \frac{\partial \zeta_i}{\partial y} \right] dR$$
$$+ \int_{\partial R} \frac{\partial}{\partial n} (\Delta^2 u_n) \zeta_i \, d\Gamma$$



(*)

$$-\int_R \left(\frac{\partial^3 u_n}{\partial x^3} \frac{\partial \xi_i}{\partial x} + \frac{\partial^3 u_n}{\partial x^2 \partial y} \frac{\partial \xi_i}{\partial y} + \frac{\partial^3 u_n}{\partial x \partial y^2} \frac{\partial \xi_i}{\partial x} + \frac{\partial^3 u_n}{\partial y^3} \frac{\partial \xi_i}{\partial y} \right) dR =$$

$$= -\int_R \left(\frac{\partial^3 u_n}{\partial x^3} + \frac{\partial^3 u_n}{\partial x \partial y^2} \right) \frac{\partial \xi_i}{\partial x} dR - \int_R \left(\frac{\partial^3 u_n}{\partial x^2 \partial y} + \frac{\partial^3 u_n}{\partial y^3} \right) \frac{\partial \xi_i}{\partial y} dR =$$

$$= -\int_R \underbrace{\Delta^2 \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial \xi_i}{\partial x}}_{\text{lemma di Green}} dR - \int_R \underbrace{\Delta^2 \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial \xi_i}{\partial y}}_{\text{lemma di Green}} dR =$$

$$= \int_R \left[\frac{\partial}{\partial x} \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial^2 \xi_i}{\partial x^2} + \frac{\partial}{\partial y} \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial^2 \xi_i}{\partial x \partial y} \right] dR + \int_R \left[\frac{\partial}{\partial x} \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial^2 \xi_i}{\partial x \partial y} + \frac{\partial}{\partial y} \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial^2 \xi_i}{\partial y^2} \right] dR$$

$$- \int_{\partial R} \frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial \xi_i}{\partial x} d\Gamma - \int_{\partial R} \frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial \xi_i}{\partial y} d\Gamma =$$

$$= \int_R \left(\frac{\partial^2 u_n}{\partial x^2} \frac{\partial^2 \xi_i}{\partial x^2} + 2 \frac{\partial^2 u_n}{\partial x \partial y} \frac{\partial^2 \xi_i}{\partial x \partial y} + \frac{\partial^2 u_n}{\partial y^2} \frac{\partial^2 \xi_i}{\partial y^2} \right) dR$$

$$- \int_{\partial R} \left[\frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial \xi_i}{\partial x} + \frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial \xi_i}{\partial y} \right] d\Gamma$$

$$- \int_R \left(\frac{\partial^3 u_n}{\partial x^3} \frac{\partial \xi_i}{\partial x} + \frac{\partial^3 u_n}{\partial x^2 \partial y} \frac{\partial \xi_i}{\partial y} + \frac{\partial^3 u_n}{\partial x \partial y^2} \frac{\partial \xi_i}{\partial x} + \frac{\partial^3 u_n}{\partial y^3} \frac{\partial \xi_i}{\partial y} \right) dR$$

$$+ \int_{\partial R} \frac{\partial}{\partial n} \left(\frac{\partial^2 u_n}{\partial x^2} + \frac{\partial^2 u_n}{\partial y^2} \right) \xi_i \cdot d\Gamma = \int_R f \xi_i \cdot dR$$

2^a applicazione:

vedi 4-3b → *

$$\int_R \left(\frac{\partial^2 u_n}{\partial x^2} \frac{\partial^2 \xi_i}{\partial x^2} + 2 \frac{\partial^2 u_n}{\partial x \partial y} \frac{\partial^2 \xi_i}{\partial x \partial y} + \frac{\partial^2 u_n}{\partial y^2} \frac{\partial^2 \xi_i}{\partial y^2} \right) dR -$$

$$- \int_{\partial R} \left[\frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial x} \right) \frac{\partial \xi_i}{\partial x} + \frac{\partial}{\partial n} \left(\frac{\partial u_n}{\partial y} \right) \frac{\partial \xi_i}{\partial y} \right] d\Gamma +$$

$$+ \int_{\partial R} \frac{\partial}{\partial n} \left(\frac{\partial^2 u_n}{\partial x^2} + \frac{\partial^2 u_n}{\partial y^2} \right) \xi_i \cdot d\Gamma = \int_R f \xi_i \cdot dR$$

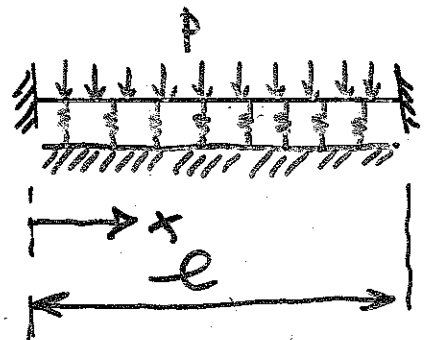
Integrali di linea = b.c. dell'eq. bi-armonica

Risultato importante: ξ_i possono ora appartenere ad uno spazio $W_2^{(2)}$ anziché $W_2^{(4)}$.

Altro esempio: trave prismatica appoggiata su fondazione elastica:

$$EI \frac{d^4 u}{dx^4} + ku = p$$

$$u(0) = u(l) = \frac{du}{dx} \Big|_{x=0} = 0$$



Integrale variazionale di Galerkin:

$$\int_0^l (EI \frac{d^4 u_n}{dx^4} + k u_n - p) \xi_i dx = 0$$

Integrando per parti:

$$-\int_0^l EI \frac{d^3 u_n}{dx^3} \frac{d \xi_i}{dx} dx + \int_0^l (k u_n - p) \xi_i dx + EI \frac{d^3 u_n}{dx^3} \xi_i \Big|_0^l = 0$$

Integrando ancora:

$$\int_0^l EI \frac{d^2 u_n}{dx^2} \frac{d^2 \xi_i}{dx^2} dx + \int_0^l (k u_n - p) \xi_i dx + EI \frac{d^3 u_n}{dx^3} \xi_i \Big|_0^l - EI \frac{d^2 u_n}{dx^2} \frac{d \xi_i}{dx} \Big|_0^l = 0$$

Resta:

$$\int_0^l EI \frac{d^2 u_n}{dx^2} \frac{d^2 \xi_i}{dx^2} dx + \int_0^l k u_n \xi_i dx = \int_0^l p \xi_i dx$$

Le funzioni base ξ_i possono appartenere a $\mathcal{K}_{(4)}^2$.